# Correlators

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### Introduction

- \* I will explain how correlator products (visibilities) are formed
- $\ast$  I will explain the fringe fit process and creation of total delays
- \* I cover several Digital Signal Processing topics
- \* This is a very mathematical subject
  - Some calculation details are in the appendix
  - $\circ\,$  Several signal processing concepts are explained along the way
  - Slow me down and ask questions as necessary!



- \* Delay model calculation already covered
- \* Making use of total delay will be covered in later talks

### Why learn about correlators?

- \* Understand interferometry data products
- \* Design interferometric experiments properly
- \* Implement or improve upon a correlator
- \* To operate a correlator
- \* To achieve an enhanced state of enlightenment

- \* Radio antennas/receivers measure electric field vectors
- \* These are handed to the correlator as voltage time series
- \* Here we are concerned with cross correlations of these
- \* 2 (or more) antennas and a correlator form a radio interferometer



- \* Definition
- \* Correlation of functions
- \* Correlation of sampled data
- \* Noise and sensitivity
- \* The complex-valued visibility

#### **Formal definition**

Any implementation the cross-correlation function,

$$C_{ij}(\tau) = \operatorname{Corr}[v_i, v_j] = \langle v_i(t)v_j(t+\tau) \rangle$$

given two real-valued functions,  $v_i(t)$  and  $v_j(t)$ .

#### **Colloquial definition**

The device that calculates the above for a VLBI (or other astronomical) observation across 2 or more antennas, each with 1 or 2 polarization components, 1 or more spectral windows with use of delay model functions  $\tau_{ij}(t)$  appropriate for the source being studied. The VLBI correlator may also extract pulse cal tones and apply certain calibration to data.

#### **Schematic**



- \* The calculatd value,  $C_{ij}(\tau)$ , is a statistical quantity
  - Must average over many (independent) samples to be meaningful
  - $\circ~$  For a bandwidth of  $\Delta\nu,$  one independent sample every  $\Delta t=1/2\Delta\nu.$
- \* Calculation is generally explicitly time-bounded
- \* Usually is computed on uniformly sampled data:

$$C_{ij}[k] = \frac{1}{N} \sum_{l=1}^{N} v_i[l] v_j[l+k]$$

with integer k and l

 $\ast \, \, k \, \, {\rm or} \, \, \tau \, \, {\rm is} \, \, {\rm called} \, \, {\rm the} \, \, {\it lag}$ 

# Example 1

- \* Use signals  $v_1(t) = \sin 2\pi\nu t$  and  $v_2(t) = \cos 2\pi\nu t$ .
- \* Take limiting case as time range extends infinitely.

$$C_{12}(\tau) = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} \cos 2\pi\nu t \sin 2\pi\nu (t+\tau) dt$$
$$= -\frac{1}{2} \sin 2\pi\nu \tau$$

- \* Narrow-band signals have large support over  $\tau$ .
- \* Sums of pure tones (as here) have support even as  $| au| o \infty$ .
- \* See appendix for detailed derivation.

#### Normalized correlation coefficient

\* Often one is interested in a normalized value (independent of scale)

$$\Gamma_{ij}(\tau) = \frac{\langle v_i(t)v_j(t+\tau)\rangle}{\sqrt{\langle v_i(t)^2\rangle \langle v_j(t+\tau)^2\rangle}}$$

- \* The denominator is the geometric mean of the two signals' autocorrelations
- \*  $\Gamma_{ij}$  is a measure of how similar the two signals are  $\circ \ \Gamma_{ij}(\tau) = \pm 1$  if and only if  $v_i(t) \propto \pm v_j(t+\tau)$ .  $\circ \$ Otherwise  $|\Gamma_{ij}(\tau)| < 1$
- \* For  $v_1(t) = \sin 2\pi\nu t$  and  $v_2(t) = \cos 2\pi\nu t$ :

$$\Gamma_{ij}(\tau) = -\sin 2\pi\nu\tau$$

\* Thus the cosine function is the same as the sine function with a  $n-1/4~{\rm period}$  shift.

# Example 2

- \* Cross-correlate  $v_1(t) = e^{-t^2/2}$  against  $v_2(t) = e^{-(t-t_0)^2/2}$
- \* For simplicity make use of  $v_1(t) = v_2(t+t_0)$ :

$$\Gamma_{12}(\tau) = \frac{\int_{-\infty}^{\infty} v_1(t) v_2(t+\tau) dt}{\int_{-\infty}^{\infty} v_1(t)^2 dt} \\ = e^{-(\tau-t_0)^2/4}$$

- \* Result could be predicted without grungy math:
  - Correlation of time symmetric signals is equivalent to convolution
  - Convolution of two Gaussians is a wider Gaussian (sum in quadrature)
  - $\circ~$  Signals are the same when  $\tau=t_0$
- \* More complete derivation in appendix

# Example 2 (continued)



### Correlation of sampled data

- \* Sampled data can be treated in similar manner as a continuous function
- \* Replace integrals by sums
- $\ast\,$  Assume here that sampled data streams both be uniformly sampled at same interval,  $\Delta t$
- \* Sampled signals must be band-limited with  $\Delta \nu \leq 1/2\Delta t$  (Nyquist sampling theorem)
- \* Note: sampled does not imply quantized; ignore quantization here
- \* Given  $v_i[l]$  and  $v_j[l]$ , the corresponding quantities are:

$$C_{ij}[k] = \frac{1}{N} \sum_{l=1}^{N} v_i[l] v_j[l+k]$$
  

$$\Gamma_{ij}[k] = \frac{\sum_{l=1}^{N} v_i[l] v_j[l+k]}{\sqrt{\sum_{l=1}^{N} v_i[l]^2} \sqrt{\sum_{l=1}^{N} v_j[l+k]^2}}$$

#### **Example 3: Seismology**



Image from Einstein Telescope design study document, 2011

# Sampling band-limited signal: original signal



### Capture signal every unit interval



### **Retain only samples**



18/104

### Naive signal reconstruction



19/104

# The interpolation function for $1^{st}$ Nyquist zone (sinc)



<sup>20/104</sup> 

### **Properly interpolated function**



# Comparison of original and reconstructed signals



22/104

- \* Take 2 sampled signals,  $\mathfrak{g}_1[l]$  and  $\mathfrak{g}_2[l]$ , where
  - $\circ\;$  Each  $\mathfrak{g}_i[k]$  is drawn from a zero mean, unit norm Normal distribution
  - $\circ~\langle \mathfrak{g}_i 
    angle = 0$ ,  $\left< \mathfrak{g}_i^2 \right> = 1$  (which implies  $C_{ij} = \Gamma_{ij}$ )
  - $\circ~\langle \mathfrak{g}_{j}\mathfrak{g}_{j}\rangle = \delta_{ij}$  (defines uncorrelated noise)
- \* The expectation value of the correlation function vanishes

$$C_{12}[k] = \frac{1}{N} \sum_{i=1}^{N} \mathfrak{g}_1[l] \mathfrak{g}_2[l+k] = 0$$

\* But its RMS does not

$$\sigma_{C_{12}[k]} = \frac{1}{\sqrt{N}}$$

\* This is the basis for calculating interferometer sensitivity (see appendix)

# The (optical) double slit experiment

- \* The film at the image plane of a double slit is a correlator!
- \*  $\tau = \tau_2 \tau_1$  is the path length difference
- \* Monochrome signal hits mask  $v(t) = \cos 2\pi\nu t$
- \* Signals at image:

$$v_1(t) = \cos 2\pi\nu(t-\tau_1)$$
  
 $v_2(t) = \cos 2\pi\nu(t-\tau_2)$ 

\* Intensity at image:

$$I(\tau) \propto \left\langle (v_1(t) + v_2(t))^2 \right\rangle$$
  
=  $1 + \cos 2\pi\nu\tau$ 



- \* This is an additive correlator
- \* The constant term is the total power
- \* The brightness ripples are *fringes*

#### **Correlation of quasi-monochromatic signals**

- \* As seen before cross correlation of two equal-frequency signals gives sinusoidal response with respect to  $\tau$ .
- \* Sinusoids have two free parameters: amplitude and phase.
- \* Seems silly to need more than two measurements to completely characterize correlator response.
- \* Solution: measure two lags, separated by 90 degrees of phase!

$$C_{ij}(\tau) = C_{ij}(0)\cos 2\pi\nu\tau + C_{ij}(1/4\nu)\sin 2\pi\nu\tau$$

\* For convenience, bundle into a single complex number

$$V_{ij} = C_{ij}(0) + iC_{ij}(1/4\nu)$$

\* This is proportional to the familiar visibility. And then

$$V_{ij}(\tau) = \operatorname{Re}\left(V_{ij}e^{-2\pi i\nu\tau}\right)$$

- \* The complex correlator
- \* The Hilbert transform
- \* Analytic signals
- \* Complex sampling

#### Schematic of complex correlator



 $V_{ij}(\tau) = \langle v_i(t)v_j(t+\tau) \rangle + i \langle v_i(t)\mathcal{H}[v_j](t+\tau) \rangle$ 

 $\ast\,$  Given a real-valued signal v(t), define analytic signal

$$w(t) = v(t) + i\mathcal{H}[v(t)]$$

 $\ast~$  Here  ${\cal H}$  is the Hilbert transform

$$\circ \ \mathcal{H}[v(t)] = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{v(s)}{t-s} ds$$

$$\circ \cos \rightarrow \sin$$
 and  $\sin \rightarrow -\cos$ 

- $\circ~\mathcal{H}[\mathcal{H}[v(t)]] = -v(t)$  (the operation is invertable)
- \* Analytic signals are mathematical tools
  - Allows complex multiplication
  - Simplifies Fourier transforms
  - Simplifies fringe rotation
- \* Remember:  $\operatorname{Im}(w(t))$  is not physical
- \* See https://en.wikipedia.org/wiki/Analytic\_signal for a good discussion

#### Hilbert transform example: amplitude modulated signal



#### Compute Hilbert transform (blue + i green is analytic)



#### Reconstruct envelope (sum blue & green in quadrature)



### Analytic signal properties

\* Energy content is double:

$$\int w(t)^* w(t) dt = 2 \int v(t)^2$$

\* Fourier transform has no negative frequency components:

$$\mathcal{F}[w(t)](\nu) = 2H(\nu)\mathcal{F}[v(t)](\nu)$$

where the Heaviside step function is:



#### Making complex sampled data from real sampled data

- \* Start with a real sampled signal, v[k]
- \* Can compute the sampled equivalent of an analytic signal using discrete Hilbert transform

$$\mathcal{H}(v)[k] = \begin{cases} \frac{2}{\pi} \sum_{n \text{ odd}} \frac{v[n]}{k-n} & k \text{ even} \\ \frac{2}{\pi} \sum_{n \text{ even}} \frac{v[n]}{k-n} & k \text{ odd} \end{cases}$$

- \* Resultant signal,  $v[k] + i\mathcal{H}(v)[k]$ , carries duplicate information
- \* Can drop alternate samples to define:

$$w[k] = v[2k] + i\mathcal{H}(v)[2k]$$

\* Note that sample rate simply is inverse bandwidth:  $\Delta t = 1/\Delta \nu$  $\circ$  Clock rate of digital electronics can be halved!

- \* Make use of the power of complex numbers
- \* Note use of complex conjugation: \*

$$Corr[w_i, w_j] \equiv \langle w_i^* w_j \rangle$$
  
=  $\langle v_i v_j \rangle + i \langle v_i \mathcal{H}[v_j] \rangle - i \langle \mathcal{H}[v_i] v_j \rangle + \langle \mathcal{H}[v_i] \mathcal{H}[v_j] \rangle$   
=  $2 \langle v_i v_j \rangle + 2i \langle v_i \mathcal{H}[v_j] \rangle$   
=  $2V_{ij}$ 

- \* The above equation holds for continuous or sampled signals
- \* Equality of the second and third expressions can be shown through spectral analysis

#### Spectral decomposition of signals

\* A band-limited signal can be expressed analytically as

$$w(t) = \int_0^{\Delta \nu} e^{2\pi i t \nu} \, \tilde{w}(\nu) \, d\nu$$

- \* Where  $\tilde{w}(\nu)^* \tilde{w}(\nu)$  is proportional to the spectral power density of the signal at frequency  $\nu$ .
- \* Nyquist sampling simply captures this each  $\Delta t = 1/\Delta \nu$ :

$$w[k] = \int_0^{\Delta \nu} e^{2\pi i k \Delta t \nu} \, \tilde{w}(\nu) \, d\nu$$

\* The correlation function can be expressed as integral over frequency rather than over time:

$$V_{ij}(\tau) = \int_{\nu_1}^{\nu_2} e^{2\pi i\nu\tau} \,\tilde{w}_i(\nu)^* \,\tilde{w}_j(\nu) \,d\nu$$

- \* Concept
  - $\,\circ\,$  First cross-multiply and accumulate (X)
  - $\circ$  Then Fourier transform (F)
- \* Spectral response
- \* Realization of lag correlators in practice
- \* Examples of lag correlators
- \* Fourier transforming a time series leads to its spectrum
- $* V_{ij}( au)$  is a time series in au
- \* What if we discrete Fourier transform it?
- $\ast\,$  Assume n lags, each spaced by the sample rate,  $\Delta t$
- \*  $V_{ij}$  is complex-valued, so total bandwidth is  $\Delta 
  u = 1/\Delta t$

$$\tilde{V}_{ij}[l] \equiv \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} e^{-2\pi i k l/n} V_{ij}[k]$$
$$= \int_0^{\Delta \nu} A_n \left(\frac{l}{n} - \nu \Delta t\right) \tilde{w}_i(\nu)^* \tilde{w}_j(\nu) d\nu$$

- \* Where did this come from? See appendix for details.
- \* What does it mean?

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### Interpretation

\* Lag correlator response is equivalent to complex correlator with additional factor

$$A_n(x) = \sin n\pi x / \sin \pi x$$

- \* The function  $A_n$  serves as a filter response
- \* Each output channel, l, has its own filter, shifted by  $\Delta 
  u/n$



#### Hanning smoothing

- \* Damp oscillitory spectra by smoothing with kernel  $(\frac{1}{4}, \frac{1}{2}, \frac{1}{4})$
- \* Causes wider but much more contained spectral response
  - $\circ~$  Can throw out every other channel without loss of information
- \* Effective in reducing impact of RFI



# Comparison: with and without Hanning smoothing

- \* Two spectra with same number of channels
- \* Second one improved but comes at higher computational cost
- \* Implications for RFI immunity?



#### Comparison: with and without channel averaging

- \* Two spectra with same number of channels
- \* Simple channel averaging does rather poor job (Gibb's effect)



- For symmetry sake want equal number of positive and negative lags
   This is actually important when considering *closure quantities*
- \* We haven't yet discussed fractional sample correction
  - $\circ~$  This allows calculation at  $\tau \neq n \Delta t$
- \* Thus an odd number of lags is natural to consider
- \* All results generalize to even and odd numbers

- \* Conceptually same as complex lag correlators
- \* Need twice as many real lags for same response
  - Each lag is half as long (duration of a real sample rather than analytic complex sample)
- $\ast$  Half as many multipliers needed, but they run at twice the rate
- \* Use real-to-complex Fourier transform
- \* Spectral expansion of signals uses sines and cosines
- \* Both real and complex lag correlators used in practice

# Schematic of (real) lag correlator



\* Note: FFT usually performed in software even, on hardware correlators

## **Examples of lag correlators**

- \* Mark4 (JIVE, Haystack, WACO, Bonn)
- \* 1997-present (mostly retired)



- \* Old VLA Correlator
- \* 1980-2008



- \* Filter banks
- \* Concept
  - First Fourier transform (F)
  - $\circ$  Then cross-multiply and accumulate (X)
- \* Spectral response
- \* Realization of FX correlators in practice

# **FFT** filter banks

- \* FFT incoming bandwidth  $\Delta \nu$  real signal in blocks of 2n $\circ$  Shown below
- \* or FFT incoming bandwidth  $\Delta 
  u$  complex signal in blocks of n
- \* Produce n complex time series, each with bandwidth  $\Delta 
  u/n$



# FFT filter bank frequency response

\* Starting from a complex sampled signal, the filter bank output is:

$$\begin{split} \tilde{w}[l] &= \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} e^{-2\pi i k l/n} w[k] \\ &= \int_{0}^{\Delta \nu} \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} e^{2\pi i k (\Delta t \nu - l/n)} \tilde{w}(\nu) d\nu \\ &= \int_{0}^{\Delta \nu} A_n \left(\frac{l}{n} - \nu \Delta t\right) \tilde{w}(\nu) d\nu \end{split}$$

- \* Note symmetric summation; through universal relabeling of samples by 1/2 sample, an even number of samples can be accomodated.
  - Not possible in lag case because the parameter was the lag itself.
  - $\circ~$  The process is equivalent to a shifted FFT



\* The visibility is computed as

$$\begin{split} \tilde{V}_{ij}[l] &= \langle \tilde{w}_i[l]^* \, \tilde{w}_j[l] \rangle \\ &= \int_0^{\Delta \nu} \left[ A_n \left( \frac{l}{n} - \nu \Delta t \right) \right]^2 \, \tilde{w}_i(\nu)^* \, \tilde{w}_j(\nu) \, d\nu \end{split}$$

\* Similar to lag correlator response but with extra factor of  $A_n()$   $\circ$  Each filterbank contributes one factor

# FX correlator frequency response

 $\ast\,$  Each channel's response is similar to that of the  ${\rm sinc}^2$  function:

$$A_n(x)^2 = \left(\sin n\pi x / \sin \pi x\right)^2$$

\* Generally better than lag corretor output but worse than Hanning smoothed lag correlator output.



#### Comparison: with and without channel averaging

- \* Two spectra with same number of channels
- \* Simple channel averaging
- \* Neighboring channels fairly well isolated
- \* Peak sidelobes still rather high



# **Examples of FX correlators**

\* VLBA hardware correlator\* 1992-2009



- \* Most software correlators (e.g., DiFX and SFXC)
- \* Not tied to particular hardware
- \* DiFX can run on a Raspberry Pi!



# What information is needed to correlate?

Time for you to brainstorm...

# What information is needed to correlate?

- \* Start and stop times
- \* Frequencies of observation + bandwidth
- \* Location of the data
- \* Format of the data
- \* Location of antennas
- \* Coordinates of the source
- \* Clock offsets
- \* Correlator parameters: time and spectral resolution

# Part 5: Fractional sample delay and fringe rotation

- \* Effect of delay error
- \* Fractional sample delay compensation
- $\ast$  Fringe rotation

- \* Assume a broadband signal of uniform spectral density  $|\tilde{w}(\nu)| = 1$
- \* Look at auto-correlation with a time lag of au
- \* Consider one correlator channel with ideal spectral response between  $\nu_1$  and  $\nu_2$ .

$$C_{ii}(\tau) = \frac{1}{\nu_2 - \nu_1} \int_{\nu_1}^{\nu_2} d\nu \, e^{-2\pi i \tau \nu} \tilde{w}_i(\nu)^* \tilde{w}_i(\nu)$$
  
$$= \frac{1}{\nu_2 - \nu_1} \int_{\nu_1}^{\nu_2} d\nu \, e^{-2\pi i \tau \nu}$$
  
$$= e^{-2\pi i \tau \nu_0} \frac{\sin \pi \tau \Delta \nu}{\pi \tau \Delta \nu}$$

\* Where  $u_0 = rac{1}{2} \left( 
u_1 + 
u_2 
ight)$  is the channel center frequency

\* And  $\Delta 
u = 
u_2 - 
u_1$  is the channel bandwidth

- \* There are two effects:
  - $\circ~$  There is a phase shift of  $2\pi\tau\nu_0$
  - $\circ\,$  There is an amplitude reduction (decorrelation) by amount

$$\frac{\sin \pi \tau \Delta \nu}{\pi \tau \Delta \nu}$$

- \* The phase error is correctable
- \* The amplitude can be restored
  - But decorrelation (loss of SNR) is permanent
  - $\circ~$  This is devistating unless  $\tau \ll 1/\Delta \nu$
- \* Remember these effects for when we discuss fringe fitting...

- $*\,$  As one observes an astronomical source, the correlator delay model,  $\tau$  must change as the source moves across the sky.
- \* Source motion is smooth with time.
- \* Bulk delay is compensated by choosing which samples to correlate.
- \* Each incoming datastream can be offset from integer sample by as much as  $\pm \frac{1}{2}$  of a sample.
- \* Compensation is handled differently on different correlator architectures.
- \* Spectral line (multi-channel) correlators simplify life:  $\Delta t \ll 1/\Delta \nu$  in most cases so effective delay error is reduced.

- \* Essentially the time-dependent fractional sample compensation
- \* Various possible places to implement:
  - At end of each visibility spectrum calculation (as phase gradient)
  - During accumulation, after each FFT (as phase gradient; FX-only)
  - In time domain, directly on each sample (sample phase rotation)
- \* Magnitude depends on frequency, not bandwidth!
- \* Remember! Want to keep phase change well under 1 radian over any averaging period.

# Post-integration fringe rotation

- \* The least costly (in terms of operations)
- \* Phase applied to visibility spectrum (part of fractional sample corr.)

$$V_{12}(\tau,\nu) = e^{-2\pi i\nu\Delta\tau} V_{12}([\tau],\nu)$$

- $\ast\,$  Where  $[\tau]$  is the delay corresponding to the nearest integer number of samples, and
- \*  $\Delta \tau = \tau [\tau]$  is the fractional sample being compensated.
- \*~ This is valid when  $T \dot{\tau} \nu \ll 1$
- \* Example:  $b=1~{\rm km}$  equatorial baseline at  $\nu=1~{\rm GHz}$  at zenith passage
  - $\circ~$  Phase as function of time:  $\phi(t)=2\pi\nu b\sin(2\pi t/86400)/c$
  - $\circ~$  The fringe rate,  $\dot{\phi}(t)=4\pi^2\nu b\cos(2\pi t/86400)/(86400c),$  peaks at 1.5 rad/sec.
  - $\circ~$  Thus post-integration fringe rotation is valid for  $T\ll 0.6~{\rm sec}$
- \* Often done on sub-integration basis.

- \* With higher fringe rates, fringe rotation must be done on shorter timescales.
- \* FX correlators expose the spectrum after each FFT.
- \* Typical continuum correlator output has frequency resolution of 0.25 MHz, implying FFT timescales of  $4\mu$ s.
- $*\,$  On a 8611 km baseline (longest VLBA), this is OK for  $\nu\ll 20$  GHz.
- \* Use with care on continent-scale VLBI arrays!

- \* This is the most common form of fringe rotation used by VLBI
- \* Simply multiply each sample by  $e^{2\pi i \nu_0 \Delta \tau}$  before correlating
- \* This makes for a complex-valued signal
  - But it is not an analytic signal!
- \* Note! This technique only works well for small fractional bandwidths
  - Same phase applied to all frequencies
  - $\circ\,$  Results in decorrelation near band edges by  ${\rm sinc}\left(\pi\frac{\Delta\nu}{\nu_0}\right)$
  - Worst cast at VLBA: 128 MHz BW centered around 1.28 GHz
    - $\blacktriangleright~1.6\%$  decorrelation at band edge
    - $\blacktriangleright$  0.5% decorrelation averaged over band
    - This is still generally acceptable

 $\circ~$  Decorrelation grows as  $\left(\frac{\Delta\nu}{\nu_0}\right)^2$ 

### Part 6: The DiFX correlator

# What DiFX (Distributed FX) does

- \* Decode incoming data
- \* Select data (coarse time delay)
- \* Fringe rotate
- \* Fourier transform
- \* Select sideband
- \* Apply fractional delay correction
- \* Cross-multiply
- \* Short-term accumulate
- \* Long-term accumulate
- \* Write visibility to disk
- \* You will run DiFX at the demo tomorrow

- $\ast$  Quantization
- \* Pulsar gating
- \* Other correlator functionality
- \* Design trade-offs

# **Quantization noise**





- \* Values determined so as to minimize quantization noise
- \* Quantization efficiency  $\eta_Q = 64\%$
- \* Effective number of bits, ENOB = 1



- \* Optimal values:  $v_0 = 0.96\sigma$ ;  $R = 3.3359 \longrightarrow \eta_Q = 88\%$
- \* ENOB = 1.92

\*  $\alpha = 0.4780\,\sigma$  determined so as to minimize quantization noise

Note: Different conventions for the codes exist (e.g., Mark5B, VDIF)

## Quantization noise distribution



# **Quantization correction**

- At low correlation, quantization decreases correlation
- Quantization causes predictable non-linearity at high correlation
- \* Linear correction is easy; full correction is complicated ...



#### Quantization in the spectral domain

- \* 1-bit quantization is extreme case of harmonic distortion
- \* Power gets scattered into harmonics
- \* Oversampling allows partial discrimination of unwanted harmonics
  - Increases signal to noise
  - $\circ~$  At a substantial data transmission cost
  - $\circ~$  Very quickly diminishing returns; better to use more bits


# **Pulsar gating**

- \* Pulsars emit regular pulses with small duty cycle
- \* Period in range 1 ms to 8 s; usually  $\Delta t \ll P_{
  m pulsar} < T$
- \* Blanking during off-pulse improves sensitivity
- \* Propagation delay is frequency dependent: best done on FX architecture



- \* Pulse cal extraction
- \* Switched power extraction
- \* Data weights
- \* Multiple phase centers
- \* Spectral zooming
- \* Band matching
- \* Overlapped FFTs

- \* Hardware advantages
  - $\circ~$  Can be 10-100  $\times~$  faster
  - $\circ~$  Can be 10-100  $\times~$  more power efficient
  - Predictable operations once commissioned (usually)
  - Guaranteed real-time performance
- \* Software advantages
  - Short development timescales
  - COTS Hardware: cost effective
  - Generally more flexible
  - $\circ~$  Extensible, even after deployed
- \* GPU-based correlators straddle the two
  - Higher compute density than CPUs
  - Less flexibility than CPUs
  - $\circ~$  More difficult development than CPUs

## Trade-offs: lag or FX architecture?

## \* Lag (XF) advantages

- Can implement weights more precisely
- $\circ~$  Individual operations can be performed with small word sizes
- $\circ~$  Access to uncorrupted lag spectrum
  - Improved quantization correction
- \* FX advantages
  - Many fewer operations (increasingly so with larger spectra)
  - Improved native spectral response
  - $\circ~$  Access to frequency domain on short timescales
    - Zoom bands and band-matching
    - More effective pulsar gating
    - ► Sub-integration RFI characterization

# Spectral response and delay window duality



- \* Related by Fourier transform
- \* Must take into consideration when calculating fringe SNR!

- \* Example: Jansky VLA's WIDAR correlator
- \* 2008-present
- \* "Filter-bank XF" architecture
- \* Filterbank forms complex-valued sub-bands
- \* Each sub-band feeds a complex lag correlator





Left: WIDAR during construction

Right: WIDAR baseline board

- \* 1- or 2-bit quantization?
- \* What spectral resolution (or number of lags) is needed?
- \* What time resolution is needed?
- \* Do I need to generate all polarization products?

- \* 1-bit sampling
  - $\circ~$  Quantization efficiency  $\eta_Q=0.636$
  - Simplest to implement
- \* 2-bit sampling
  - $\circ \ \eta_Q = 0.882$
  - $\circ \ \eta_Q/\sqrt{2} = 0.624$
  - $\circ~$  Slightly lower sensitivity at fixed bitrate
  - $\circ~$  Quantization correction more linear
    - Improved performance in RFI environment
- \* Are there better quantization schemes? Yes...
- \* Why can 1- and 2-bit quantization work?
  - $\circ~$  Absolute amplitude restored with total power measurements
  - $\circ~$  The statistics of correlation are what matter

- \* Must obey frequency-time resolution product  $\delta \nu \Delta t > 1$ 
  - Otherwise you are not correlating independent samples
  - $\circ~$  Very rarely is this a limitation
- \* At low frequency RFI excision is better with more channels; can throw out affected data
  - $\circ~$  Usually this means  $\delta\nu\ll\Delta\nu$
- $*\,$  To accommodate typical clock uncertainties, open up the delay window to at least  $\pm 2\mu {\rm s}$ 
  - $\circ~$  Implies spectral resolution of 0.25 MHz or better
- \* Number of (real) lags to accomplish is simply  $2\delta 
  u/\Delta 
  u$

- $\ast\,$  Time resolution,  $\Delta t=$  accumulation period
- \*  $\Delta t$  should be smaller than timechange of correlator statistics
  - $\circ~$  Atmospheric / ionospheric pathlength changes
  - $\circ\,$  Delay model not accurate (antenna or source position error)
    - ► Residual rates usually measured in mHz or 10s of mHz
  - RFI environment
  - $\circ~$  Source structure or spectrum change (e.g., pulsars)
- \* For most cm-wave VLBI, including most geodetic processing 1 or 2 seconds is fine
- \* For mm-wave VLBI and space VLBI, smaller number is usually needed

## Should I correlate all polarization products?

- \* It depends
- \* Generally no harm in doing so, but increases computational load and output data size
- \* Many times not possible (e.g., only one polarization recorded)
- \* Are polarizations linear?
  - Probably; polarization basis rotates on sky differently at each antenna
  - $\circ~$  Not required for array of equatorial mounted antennas
- \* Are there mixed linear and circular systems?
  - $\circ\,$  Yes, otherwise you will reduce your sensitivity on some baselines

- \* The most primitive analysis step after correlation
- \* Data: time series of visibility spectra
  - For a single source
  - $\circ~$  From within a single common spectral window
  - $\circ~$  For a particular polarization product
  - $\circ~$  Over a short enough time to prevent atmospheric decorrelation
  - $\circ~$  Over a long enough time to achieve sensitivity requirements
- \* Remember: consequence of incorrect delay model:

$$\Delta\phi(\nu) = 2\pi\tau\nu$$

where  $\boldsymbol{\tau}$  here is interpreted as the delay error

- $\ast\,$  Residual phase,  $\Delta\phi,$  is proportional to frequency
- \* Delay error cause: antenna/source positions, clock error, atmosphere

## Delays

\* Phase delay

$$D_{\phi} = -\frac{1}{2\pi} \frac{\phi}{\nu}$$

\* Group delay

$$D_G = -\frac{1}{2\pi} \frac{d\phi}{d\nu}$$

- \* In non-dispersive media (e.g., vacuum) these are equal
- \* In dispersive media (e.g., ionosphere) they differ
- \* Group delay is the direct observable of VLBI observations
- $\ast\,$  Phase delay suffers from  $2\pi$  ambiguities

#### Rates

\* Phase rate

$$R_{\phi} = \frac{1}{2\pi} \frac{d\phi}{dt}$$

\* Delay rate

$$R_D = \frac{1}{2\pi\nu} \frac{d\phi}{dt}$$

- \* The two quantities are easily and unambiguously convertable
- \* Fringe fitting naturally produces phase rate

## **Example data**



Frequency (spans 256 MHz)

- $\ast\,$  Two adjacent 128 MHz bands on HN-LA baseline at 4.8 GHz
- $\ast\,$  By eye: approx 5.5 turns of phase across 256 MHz
- $\rightarrow D_G \sim 5.5/256 \text{MHz} = 21.5 \text{ns}$ 
  - \* AIPS FRING result:
- $\longrightarrow D_G = 21.7 \text{ns} \ R_{\phi} = -0.6 \text{mHz} \ R_D = -1.3 \times 10^{-13} \text{sec/sec}$

- \* Observe sufficiently bright, sufficiently point-like source
- \* Antennas' bandpass response relatively phase flat
- \* During solution interval delay error evolves linearly in time
- \* Then phase can be expressed as:

$$\phi(\nu, t) = \phi_0 - 2\pi D_G \nu + 2\pi R_\phi t$$

- \* Where  $\phi_0$  is the phase at the reference time and frequency
- \* Nice linear equation with 3 parameters, right?
- $\ast\,$  Not quite: phase only measured modulo  $2\pi\,$

- \* 2D FFT on  $\phi(\nu,t)$ 
  - Maybe oversampled
- \* Identify peak valued pixel
- \* Centroid the peak  $\longrightarrow (D_G, R_{\phi}, \phi_0)$  estimate
- \* Subtract  $\phi_0 2\pi D_G \nu + 2\pi R_{\phi} t$  estimate from phases
  - $\circ\,$  Phases now all close to zero;  $2\pi$  ambiguities less important
- \* Perform least-squares fit to these residuals
- \* Add fit values to estimates

# Example data (corrected)



Frequency (spans 256 MHz)

- \* Remaining phase ripple due to antenna bandpass
- \* Note increased phase noise near band edges

- \* No FFT peaks  $\longrightarrow$  no solution
- $* \ \, \mathsf{Multiple} \ \, \mathsf{FFT} \ \, \mathsf{peaks} \longrightarrow$
- \* Aliased fringes

# Appendices

- \* Trigonometric identities
- \* Symmetric power series sum
- \* Correlation of cosine and sine functions
- \* Correlation of Gaussian pulses
- \* Correlation of signals with noise
- \* Interferometer sensitivity

# **Trigonometric identities**

$$\cos(x \pm y) = \cos x \cos y \mp \sin x \sin y$$
  

$$\sin(x \pm y) = \sin x \cos y \pm \cos x \sin y$$
  

$$\cos x \cos y = \frac{1}{2} [\cos(x - y) + \cos(x + y)]$$
  

$$\cos x \sin y = \frac{1}{2} [\sin(x - y) + \sin(x + y)]$$
  

$$\sin x \sin y = \frac{1}{2} [\cos(x - y) - \cos(x - y)]$$
  

$$e^{ix} = \cos x + i \sin x$$
  

$$\cos x = \frac{e^{ix} + e^{-ix}}{2}$$
  

$$\sin x = \frac{e^{ix} - e^{-ix}}{2i}$$

# Symmetric power series sum

$$\begin{aligned} A_{2m+1}(x) &= \sum_{k=-m}^{m} e^{-i2\pi xk} \\ &= \sum_{k=-m}^{\infty} e^{-i2\pi xk} - \sum_{k=m+1}^{\infty} e^{-i2\pi xk} \\ &= \left( e^{i(2m)\pi x} - e^{-i(2m-2)\pi x} \right) \sum_{k=0}^{\infty} e^{-i2\pi xk} \\ &= \left( e^{i(2m)\pi x} - e^{-i(2m-2)\pi x} \right) \frac{1}{1 - e^{-i2\pi xm}} \\ &= \frac{e^{i(2m+1)\pi x} - e^{-i(2m+1)\pi x}}{e^{i\pi xm} - e^{-i\pi xm}} \\ &= \frac{\sin(2m+1)\pi x}{\sin\pi x} \longrightarrow A_n(x) = \frac{\sin n\pi x}{\sin\pi x} \end{aligned}$$
  
in the limit that  $n \to \infty$ ,  $A_n(x) \to \delta(x)$ ,  $\frac{A_n(x/n)}{n} \to \sin 2\pi x$ .

Note:

## The lag correlator in detail

$$\begin{split} \tilde{V}_{12}[l] &\equiv \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} e^{-2\pi i k l/n} V_{12}[k] \\ &= \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} \frac{1}{N} \sum_{j=1}^{N} e^{-2\pi i k l/n} w_1[j]^* w_2[j+k] \\ &= \int_0^{\Delta \nu} d\nu_1 \int_0^{\Delta \nu} d\nu_2 \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} \frac{1}{N} \sum_{j=1}^{N} e^{-2\pi i k l/n} \\ &\times e^{2\pi i \Delta t j \nu_1} \tilde{w}_1(\nu_1)^* e^{2\pi i \Delta t (j+k) \nu_2} \tilde{w}_2(\nu_2) \end{split}$$

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## The lag correlator in detail ...

$$\begin{split} \tilde{V}_{12}[l] &= \int_{0}^{\Delta \nu} d\nu_{1} \int_{0}^{\Delta \nu} d\nu_{2} \sum_{k=-\frac{n-1}{2}}^{\frac{n-1}{2}} e^{-2\pi i k \left(\frac{l}{n} - \nu_{2} \Delta t\right)} \\ &\times \frac{1}{N} \sum_{j=1}^{N} e^{2\pi i \Delta t j (\nu_{2} - \nu_{1})} \tilde{w}_{1}(\nu_{1})^{*} \tilde{w}_{2}(\nu_{2}) \\ &\sim \int_{0}^{\Delta \nu} d\nu_{1} \int_{0}^{\Delta \nu} d\nu_{2} A_{n} \left(\frac{l}{n} - \nu_{2} \Delta t\right) \delta(\nu_{2} - \nu_{1}) \tilde{w}_{1}(\nu_{1})^{*} \tilde{w}_{2}(\nu_{2}) \\ &= \int A_{n} \left(\frac{l}{n} - \nu \Delta t\right) \tilde{w}_{1}(\nu)^{*} \tilde{w}_{2}(\nu) d\nu \end{split}$$

- \* Here  $A_n(x) = \sin n\pi x / \sin \pi x$
- $* A_n$  is related to the sinc function
- \* See previous appendix for Derivation of function  $A_n$

#### Correlation of cosine and sine functions w/ real correlator

$$C_{12}(\tau) = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} \cos 2\pi t \sin 2\pi (t+\tau) dt$$
  
$$= \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{T} \frac{1}{2} \left[ -\sin 2\pi\nu\tau + \sin 2\pi\nu(2t+\tau) \right] dt$$
  
$$= \lim_{T \to \infty} \frac{1}{4T} \left[ -t \sin 2\pi\nu\tau - \frac{1}{4\pi\nu} \cos 2\nu(2t+\tau) \right]_{-T}^{T}$$
  
$$= \lim_{T \to \infty} -\frac{1}{2} \sin 2\pi\nu\tau + \mathcal{O}\left(\frac{1}{T}\right)$$
  
$$= -\frac{1}{2} \sin 2\pi\nu\tau$$

#### **Correlation of Gaussian pulses**

$$\Gamma_{ij}(\tau) = \frac{\int_{-\infty}^{\infty} e^{-t^2/2} e^{-(t-t_0+\tau)^2/2} dt}{\int_{-\infty}^{\infty} e^{-t^2} dt}$$
  
=  $\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-(t+t_0/2-\tau/2)^2/2} e^{-(t-t_0/2+\tau/2)^2/2} dt$   
=  $\frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} e^{-t^2} e^{-(\tau-t_0)^2/2} dt$   
=  $e^{-(\tau-t_0)^2/2}$ 

Note use of Gaussian integral identity (twice):

$$\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{\pi}$$

# The interpolation function for $2^{nd}$ Nyquist zone $(\sin 2x - \sin x)/x$



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#### Correlation of signals with noise (at zero delay)

\*  $Corr[v_i, v_j]$  is *bilinear* in its signal arguments:

$$\operatorname{Corr}[\alpha \, a + \beta \, b, \gamma \, c + \delta \, d] = \alpha \, \gamma \operatorname{Corr}[a, c] \\ + \alpha \, \delta \operatorname{Corr}[a, d] \\ + \beta \, \gamma \operatorname{Corr}[b, c] \\ + \beta \, \delta \operatorname{Corr}[b, d]$$

\* A simplistic signal model for observation of a point source is

$$\begin{aligned} v_1[k] &= S[k] + N_1[k] = \sqrt{s} \, \mathfrak{g}_0[k] + \sqrt{n_1} \, \mathfrak{g}_1[k] \\ v_2[k] &= S[k] + N_2[k] = \sqrt{s} \, \mathfrak{g}_0[k] + \sqrt{n_2} \, \mathfrak{g}_2[k] \end{aligned}$$

- \* Where S[k] and both  $N_i[k]$  are all independent Gaussian noise streams.
- $* \mathfrak{g}_i$  are unit norm zero mean Gaussian streams.
- \* For convenience, s and  $n_i$  are dimensioned as powers.

## Correlation of signals with noise (at zero delay)

\* Make use of bilinearity and previous relations:

$$C_{ij}[0] = \langle SS \rangle + \langle N_1S \rangle + \langle SN_2 \rangle + \langle N_1N_2 \rangle$$
$$= \frac{1}{N} \sum_{l=1}^N S[l]^2$$
$$= s$$

\* And normalized correlation coefficient:

$$\Gamma_{ij}[0] = \frac{s}{\sqrt{s+n_1}\sqrt{s+n_2}}$$

\* In the low signal to noise limit ( $s \ll \min n_1, n_2$ )

$$\Gamma_{ij}[0] = \frac{s}{\sqrt{n_1 n_2}}$$

## Correlation of signals with noise (at zero delay)

\* Noise does not enter the expectation value of  $C_{ij}$ , but it does the uncertainty:

$$\sigma_{C_{ij}[0]} = \sqrt{\frac{2s^2 + n_1s + sn_2 + n_1n_2}{N}}$$

- \* Some messy statistics used, left as exercise to the astute reader!
- \* In the low signal to noise limit

$$C_{ij}[0] = s \pm \sqrt{\frac{n_1 n_2}{N}}$$
  

$$\Gamma_{ij}[0] = \frac{s}{\sqrt{n_1 n_2}} \pm \frac{1}{\sqrt{N}}$$

\* Exercise to reader: consider the strong signal case.

\* Previous page allows one to write down the SNR for a measurement:

$$SNR = \frac{s\sqrt{N}}{\sqrt{n_1 n_2}} = \frac{s\sqrt{N}}{\sqrt{SEFD_1 SEFD_2}}$$

\* Usually instead the sensitivity of the baseline is expressed:

$$\Delta S = \frac{\sqrt{\text{SEFD}_1 \text{SEFD}_2}}{\sqrt{N}} = \frac{\sqrt{\text{SEFD}_1 \text{SEFD}_2}}{\sqrt{2\Delta\nu T}}$$

- \* SEFD is the System Equivalent Flux Density
  - $\circ~{\rm SEFD}=T_{\rm sys}/g$  where g is antenna gain (units of K/Jy)
  - $\circ\,$  Equals the brightness (in Jy) of a source required to double antenna noise power  $(T_{\rm sys})$
  - $\circ\,$  VLBA antenna SEFD is typically 300 to 500 Jy.
- \* Additional efficiency factors may apply (e.g., quantization)

\* Some signals are zero outside a finite time range

 $\circ~$  Or diminish sufficiently fast such that  $\lim_{T \to \infty} \int_{-T}^{T} v(t)^2 dt = C$ 

- \* Time averages of cross- and auto-correlations  $\rightarrow 0$  as  $T \rightarrow \infty$
- \* In such cases one can take the limit as follows:

$$\Gamma_{ij}(\tau) = \lim_{T \to \infty} \frac{\frac{1}{2T} \int_{-T}^{T} v_1(t) v_2(t+\tau) dt}{\sqrt{\frac{1}{2T} \int_{-T}^{T} v_1(t)^2 dt} \sqrt{\frac{1}{2T} \int_{-T}^{T} v_2(t+\tau)^2 dt}}$$

$$= \lim_{T \to \infty} \frac{\int_{-T}^{T} v_1(t) v_2(t+\tau) dt}{\sqrt{\int_{-T}^{T} v_1(t)^2 dt} \sqrt{\int_{-T}^{T} v_2(t+\tau)^2 dt}}$$